



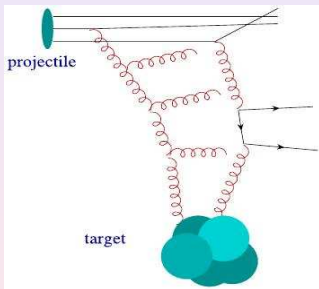
**QGSJET-III: physics, hadron production,
and applications for high energy CR studies**

**Sergey Ostapchenko
Frankfurt Institute for Advanced Studies**

**ISCRA-2019
Moscow, June 25-28, 2019**

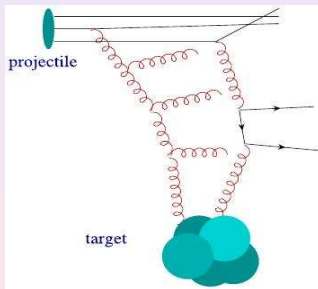
Qualitative picture for hadronic MC event generators

- QCD-inspired: **interaction mediated by parton cascades**
- **multiple scattering**
(many cascades in parallel)
- real cascades \Rightarrow particle production
- virtual cascades \Rightarrow elastic rescattering (momentum transfer)
- generally nonperturbative physics
 \Rightarrow phenomenological approaches



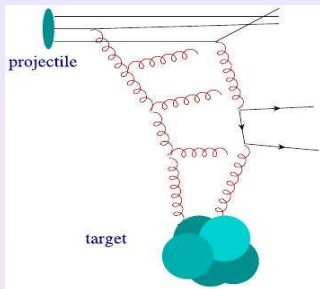
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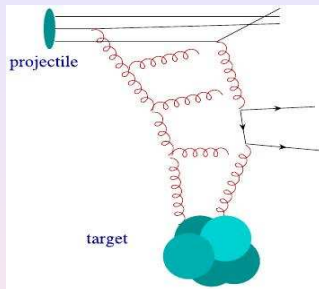


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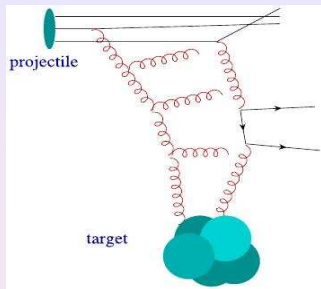


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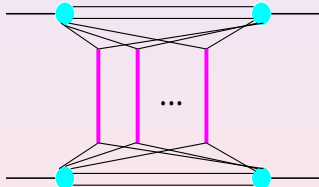
- different hadrons (nuclei) \Rightarrow different initial conditions (parton Fock states) but same mechanism
- energy-evolution of the observables (e.g. σ_{pp}^{tot}): due to a larger phase space for cascades to develop
- \Rightarrow **smooth energy-dependence for all the observables**

Soft interactions & Reggeon Field Theory (RFT)

- nonperturbative soft (small p_t) interactions:
successfully treated by RFT [Gribov, 1967]
 - Quark-Gluon String Model [Kaidalov & Ter-Martyrosian, 1982]

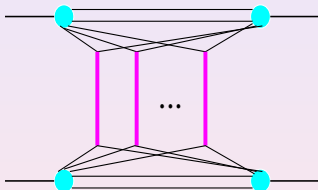
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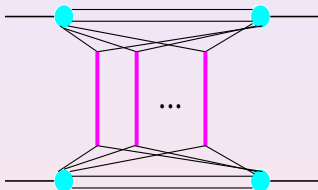


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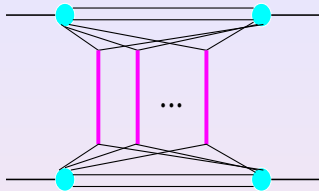
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- particle production: **hadronization of quark-gluon strings**

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Involves minimal number of adjustable parameters

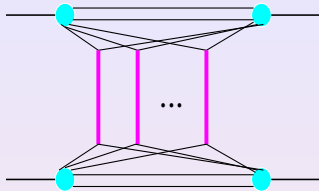
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- Pomeron intercept $\alpha_{\mathbb{P}}(0) > 1 \Rightarrow$ energy rise of parton density
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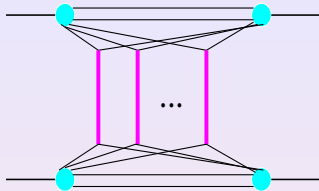
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NB: N of parameters for hadronization procedures depends on the degree of sophistication (types of secondary hadrons included, etc.)

- optionally, one may use external procedures (e.g. ones tuned to the data on e^+e^- annihilation into hadrons)

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NB: additional parameters needed to describe inelastic diffraction

- in QGSM: shower enhancement coefficients (C_{pp} , $C_{\pi p}$, C_{Kp})

- original Gribov's formulation: **assuming small parton p_t -s**
 - \Rightarrow no room for high p_t jets?

QGSJET: including hard processes in the RFT framework

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- Alternative: **treat hard processes in the RFT framework**
 - QGSJET [*Kalmykov, SO & Pavlov, 1997*]
 - neXus [*Drescher et al., 2001*]
 - EPOS [*Werner et al., 2006; Pierog et al., 2015*]

QGSJET: including hard processes in the RFT framework

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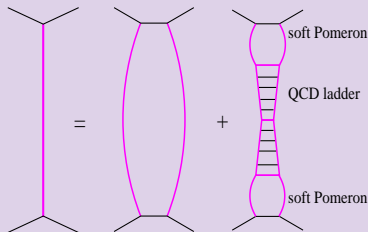
Phenomenological treatment: 'semihard Pomeron' [Drescher, Hladik, SO, Pierog & Werner, Phys. Rep. 350 (2001) 93]

- **soft Pomerons to describe soft (parts of) cascades** ($p_t^2 < Q_0^2$)
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- DGLAP for hard cascades

- taken together:
'general Pomeron'

$$\chi_{pp}^{\text{tot}}(s, b, Q_0^2) = \chi_{pp}^{\text{soft}}(s, b) + \chi_{pp}^{\text{semihard}}(s, b, Q_0^2)$$



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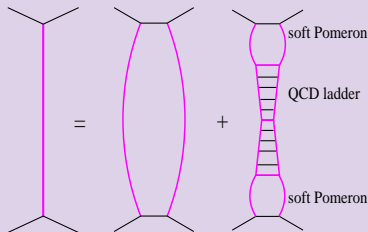
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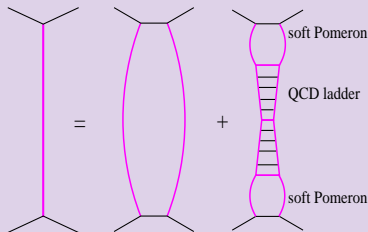
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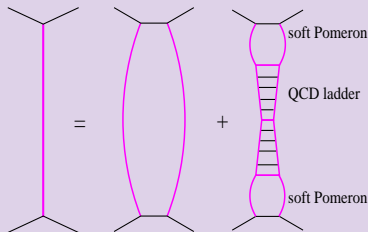
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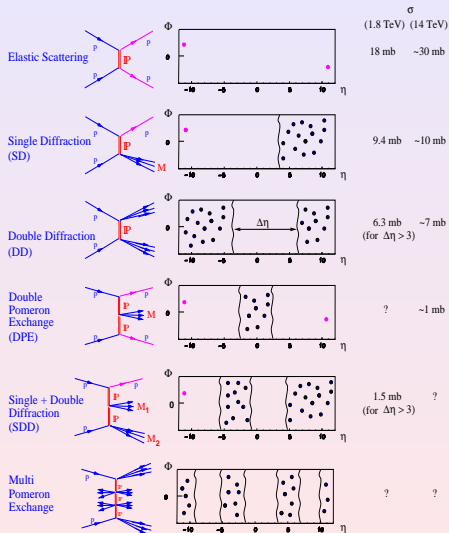
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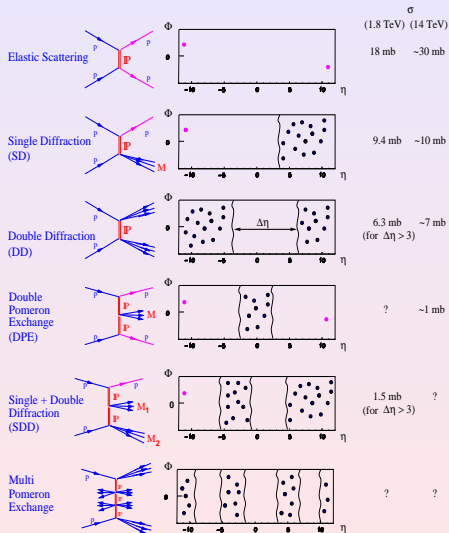
- apart from the Q_0 -cutoff, involves 2 more parameters:
to describe parton distributions in the soft Pomeron

Inelastic diffraction: Good-Walker approach and beyond



- experimentally: **formation of LRG not covered by secondaries**
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- experimentally: formation of LRG not covered by secondaries
- in many models (e.g. PYTHIA), diffraction is treated independently of ND collisions
- but: microscopically, diffractive treatment is closely related to cross sections & ND particle production (e.g. higher diffraction \Rightarrow smaller $\sigma_{pp}^{\text{inel}}$ & longer multiplicity tails)

Inelastic diffraction: Good-Walker approach and beyond

Good-Walker approach: proton is a superposition of a number of elastic scattering eigenstates: $|p\rangle = \sum_i \sqrt{C_i} |i\rangle$

$$p = \text{large light blue circle} + \text{medium dark blue circle} + \text{small dark blue circle} + \dots$$

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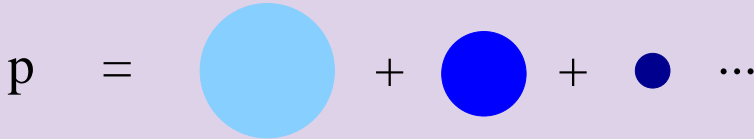
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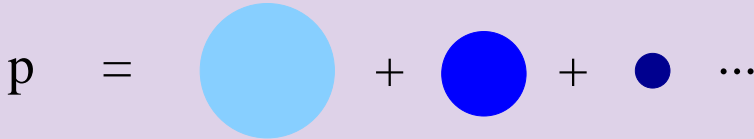


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for different combinations of such states, e.g.

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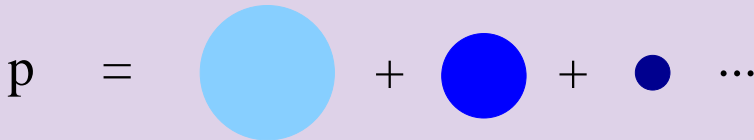
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- for each state $|i\rangle$: its own size & parton density
- should momentum sum rule be **satisfied for each state $|i\rangle$ separately**: $\sum_{l=q,\bar{q},g} \int dx x f_{l/p(i)}(x, Q^2) = 1?!$

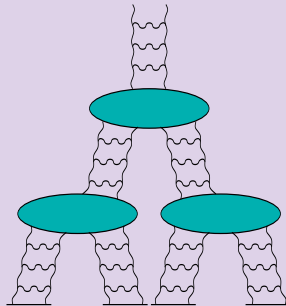
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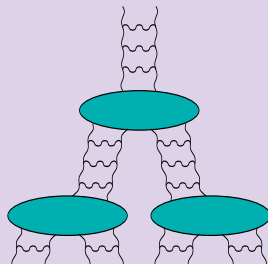
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- \Rightarrow shadowing effects
(slower rise of parton density)
- saturation: parton production compensated by fusion of partons



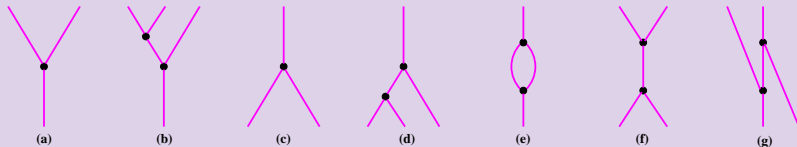
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In QGSJET-II: Pomeron-Pomeron interactions (scattering of intermediate partons off the proj./target hadrons & off each other)

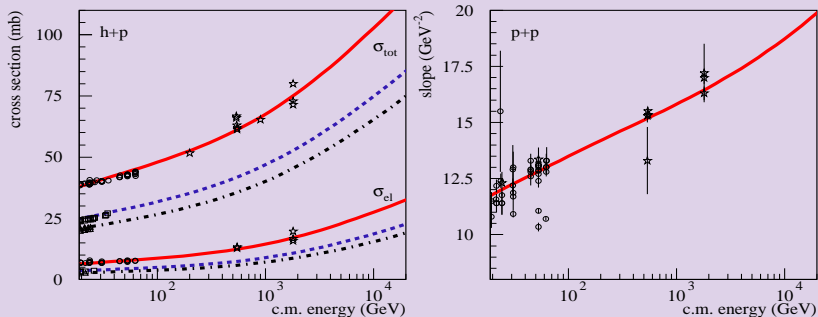


thick lines = Pomerons = 'elementary' parton cascades

- contributions resummed to all orders (sign-altering series)

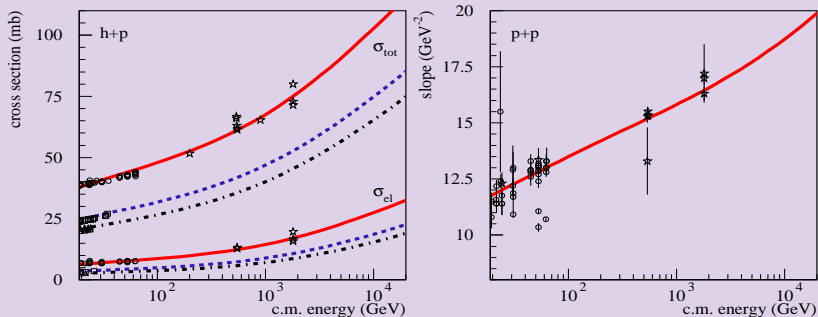
QGSJET-II-04: consistent description of $\sigma_{\text{tot/el}} & F_2$

E.g., \sqrt{s} -dependence of $\sigma_{pp/\pi p/Kp}^{\text{tot/el}}$ for realistic transverse profiles

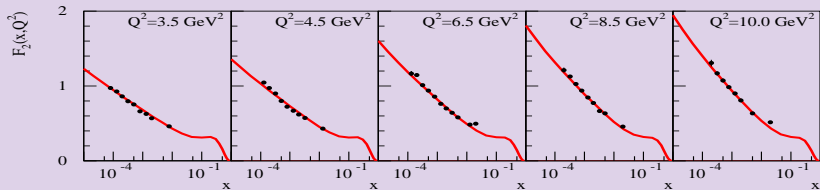


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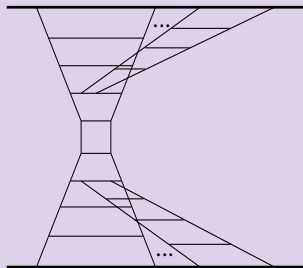


and for PDFs fitting HERA data...



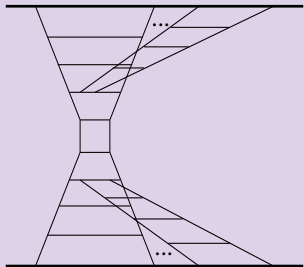
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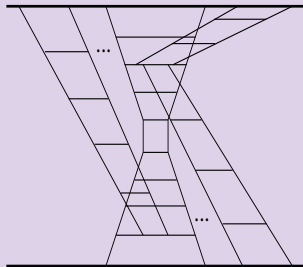
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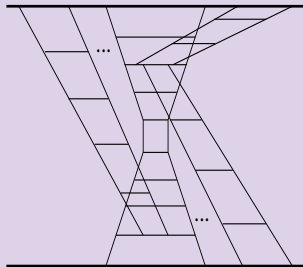
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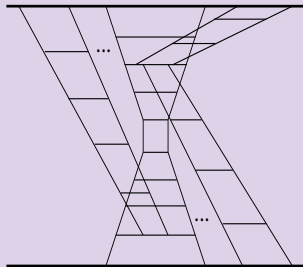
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cross sections**



This is nontrivial, not being related to parton saturation

- nonfactorizable graphs:
rescattering off the partner hadrons
- have no impact on PDFs &
inclusive particle spectra
- but: strongly damp interaction
cross sections



Many other models: energy dependent p_t -cutoff for jet production,
 $p_{t,\text{cut}} = p_{t,\text{cut}}(s)$

- is it reasonable and **what kind of physics is behind?**

QGSJET-III: treatment of higher twist (HT) effects

Any model should respect collinear factorization of pQCD

$$\begin{aligned}\sigma_{pp}^{\text{jet}}(s, p_{t,\text{cut}}) &= \sum_{I,J=q,\bar{q},g} \int_{p_t > p_{t,\text{cut}}} dp_t^2 \int dx^+ dx^- \frac{d\sigma_{IJ}^{2 \rightarrow 2}(x^+ x^- s, p_t^2)}{dp_t^2} \\ &\times f_{I/p}(x^+, M_F^2) f_{J/p}(x^-, M_F^2)\end{aligned}$$

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 - with the factorization scale $M_F^2 = p_t^2/4$, yields $p_t^{\text{cut}} \simeq 3.4$ GeV
 - but: **pQCD should work down to $Q_0 \sim 1$ GeV?!**

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- ideally, p_t -cutoff should be just a technical parameter, without a strong impact on the results
- \Rightarrow **some important perturbative mechanism seems missing**

Phenomenological approaches: higher twist (HT) effects

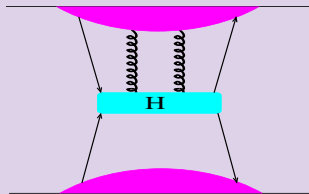
Collinear factorization: valid at leading twist (up to $1/Q^n$ terms)

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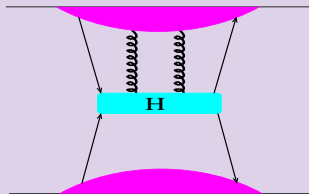
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QGSJET-III: phenomenological implementation of the mechanism

- with HT effects: **dependence on Q_0 -cutoff strongly reduced** [SO & Bleicher, 2019]
 - now: twice smaller cutoff for hard processes ($Q_0^2 = 1.5 \text{ GeV}^2$)

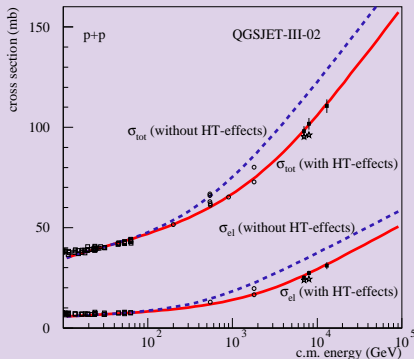
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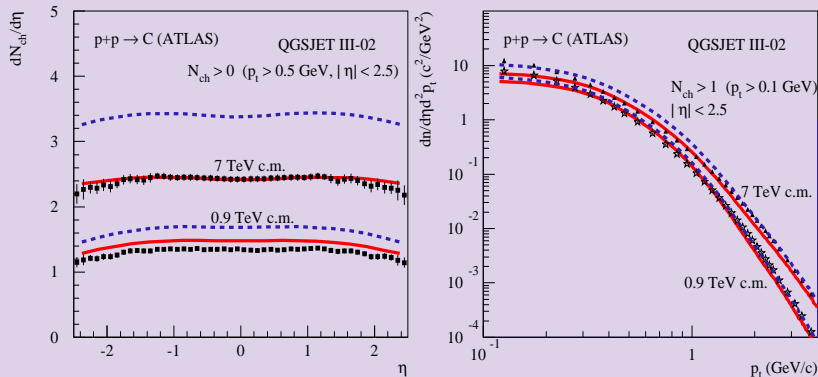
Impact on \sqrt{s} -dependence of $\sigma_{pp}^{\text{tot/el}}$

- significant corrections for total/elastic cross sections
 - start to be important already at $\sqrt{s} \sim 1 \text{ TeV}$



Phenomenological approaches: higher twist (HT) effects

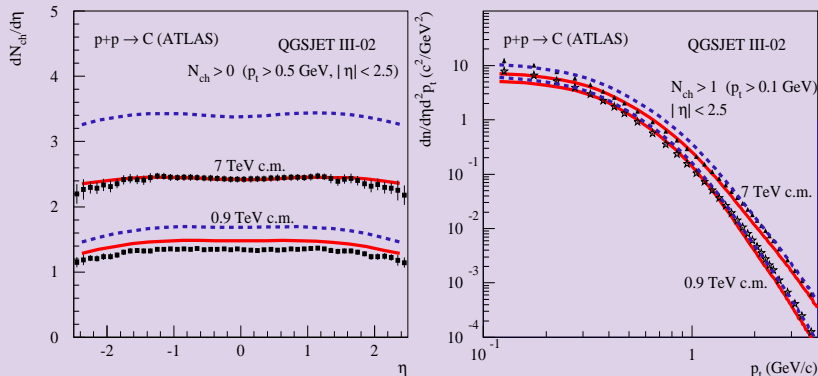
Impact on charged hadron multiplicity & p_t -spectra



- reduction of N_{ch} : stronger at higher energies
- mostly for moderately small p_t :
the effect fades away for increasing p_t ($\propto 1/p_t^2$)

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Results for air showers: preliminary and close to QGSJET-II-04

- e.g. difference for N_μ – at percent level

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 - now: twice smaller cutoff for hard processes ($Q_0^2 = 1.5 \text{ GeV}^2$)

NB: qualitatively, the approach mimics an energy dependent p_t -cutoff for jet production

- suppresses emission of jets of moderately small p_t
- has no impact on PDFs \Rightarrow not related to parton saturation

QGSJET-III: number of adjustable parameters

- basic treatment (pp , πp , Kp): 15
(soft & hard interactions; low mass diffraction)
- nonlinear effects (Pomeron-Pomeron interactions): 1
- higher twist effects: 1
- hadronization parameters: < 20

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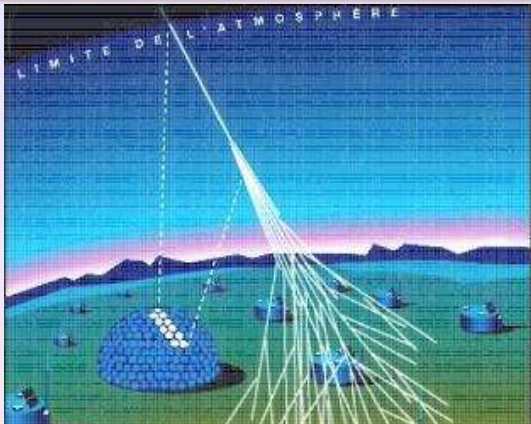
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Generally: present models of hadronic collisions
– rather involved but largely phenomenological

- \Rightarrow no wonder models differ from each other
- however: **predictions now strongly constrained by LHC data**
(using a particular model framework)

Air shower characteristics & hadronic interactions

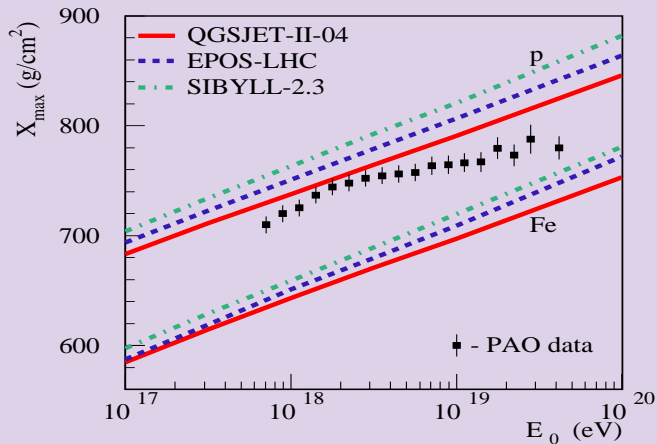


CR composition studies – most dependent on interaction models

- e.g. predictions for X_{\max} : **on the properties of the primary particle interaction** ($\sigma_{p\text{-air}}^{\text{inel}}$, $\sigma_{p\text{-air}}^{\text{diffr}}$, $K_{p\text{-air}}^{\text{inel}}$)
- predictions for muon density: **on secondary particle interactions (cascade multiplication)**; mostly on $N_{\pi\text{-air}}^{\text{ch}}$

Air shower characteristics & hadronic interactions

Why different model predictions for X_{\max} ?

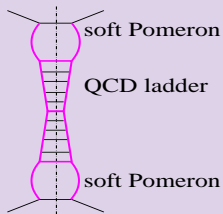


- $\sigma_{p\text{-air}}^{\text{inel}}$ – constrained by LHC studies of pp collisions
- uncertainties for $\sigma_{p\text{-air}}^{\text{diff}}$: small impact ($< 10 \text{ g/cm}^2$) [SO, 2014]
- what about $K_{p\text{-air}}^{\text{inel}}$?

Structure of constituent parton Fock states

Initial state emission (ISE) of partons doesn't stop at the Q_0 -cutoff

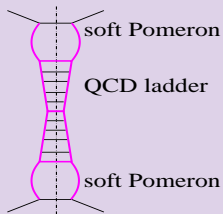
- it is extended into nonperturbative region by the soft Pomeron
- **this changes the structure of constituent parton Fock states** (represented by end-point partons in ISE)
 - in QGSJET(-II): described by Reggeon asymptotics ($\propto x^{-\alpha_{\mathbb{R}}(0)}$, $\alpha_{\mathbb{R}}(0) \simeq 0.5$)



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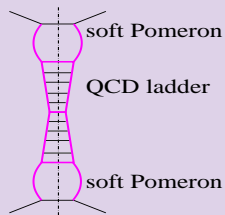
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- \Rightarrow **observables consequences:**
 - softer forward spectra (energy sharing between constituent partons)
 - **forward & central particle production - strongly correlated** (more activity in central detectors \Rightarrow larger Fock states \Rightarrow softer forward spectra)



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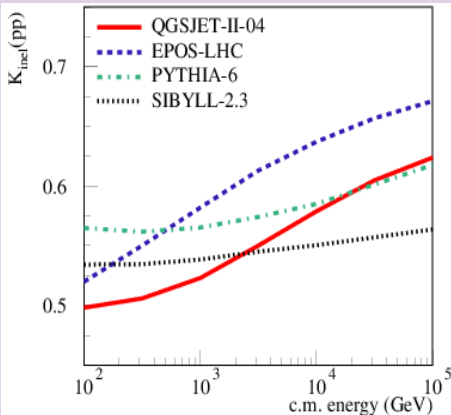


Alternative (SIBYLL & PYTHIA): no "soft preevolution"

- \Rightarrow **multiple scattering has small impact on forward spectra**
 - Feynman scaling for forward production
 - forward & central production – decoupled from each other

Structure of constituent parton Fock states

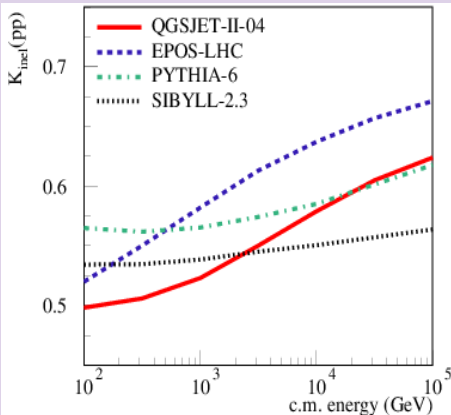
Of importance for cosmic ray studies: \sqrt{s} -dependence of K_{pp}^{inel}



- SIBYLL & PYTHIA: **weak energy dependence of the nucleon 'inelasticity'** (for increasing \sqrt{s} , mostly rise of central production)
- smaller K^{inel} \Rightarrow stronger 'leading particle' effect
- \Rightarrow slower development of CR-induced air showers

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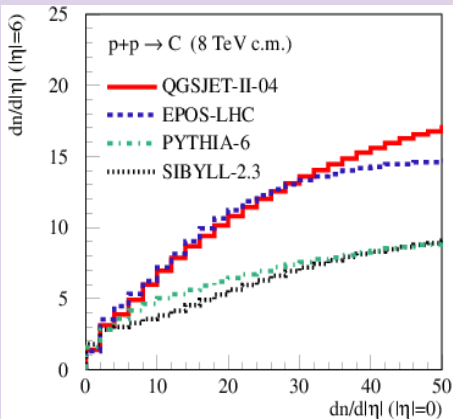
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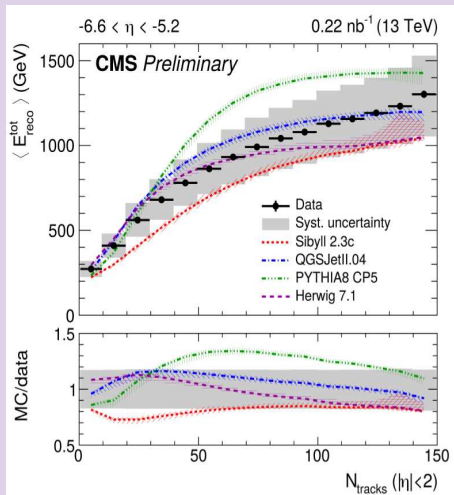
Crucial test: cross-correlation of $dN_{pp}^{\text{ch}}/d|\eta|$ at $\eta = 0$ and $\eta = 6$



- strong correlation for QGSJET-II & EPOS (apart from the tails of the N^{ch} distributions)
- twice weaker correlation for SIBYLL & PYTHIA

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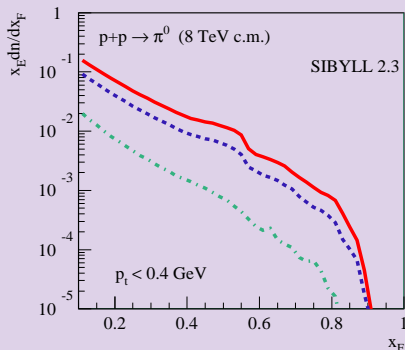
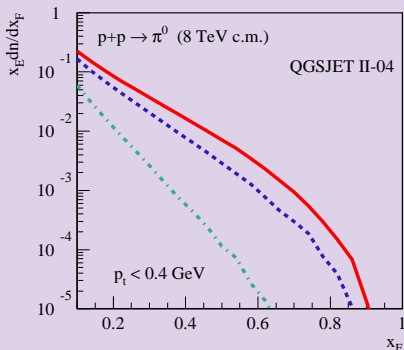
Now measured: correlation of forward energy (in CASTOR) with central activity (N of charged particle tracks) in CMS



- most important – first 3 bins ($N_{\text{tracks}} < 30$)
- very puzzling results: intermediate between QGSJET-II and SIBYLL?!
- decisive discrimination not possible?

Further discrimination: forward hadrons by LHCf & ATLAS

Forward π^0 spectra in LHCf for different ATLAS triggers
($\geq 1, 6, 20$ charged hadrons of $p_t > 0.5$ GeV & $|\eta| < 2.5$)



Compare QGSJET-II-04 (left) to SIBYLL 2.3 (right)

- enhanced multiple scattering
⇒ softer pion spectra
 - ⇒ violation of limiting fragmentation
- nearly same spectral shape for all the triggers
 - ⇒ perfect limiting fragmentation

What next?

What about other differences for EAS predictions?

- now largely dominated by model differences for pion-air (kaon-air) collisions [*SO & Bleicher, 2016*]

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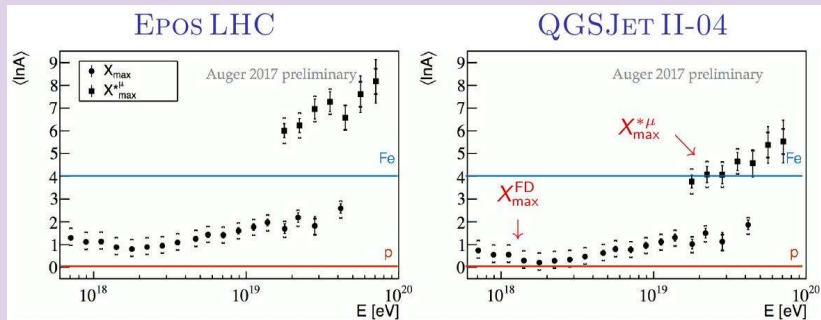
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- current indications from UHECR data (X_{\max} & X_{\max}^{μ}):
treatment of pion-air collisions may be deficient (extra slides)

Extra slides

Interpreting PAO data on X_{\max} & X_{\max}^{μ} : not self-consistent

How to change models to 'marry' X_{\max} & X_{\max}^{μ} composition-wise?

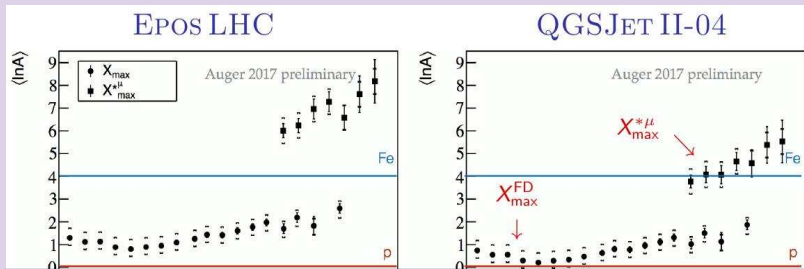


[R. Prado, ISVHECRI-2018]

- the two sets of data should overlap in terms of $\langle \ln A \rangle$
 - for $1 \leq A \leq 56!$

Interpreting PAO data on X_{\max} & X_{\max}^{μ} : not self-consistent

How to change models to 'marry' X_{\max} & X_{\max}^{μ} composition-wise?



Ancient Greek wisdom may help...



- change a model to modify X_{\max} prediction:
 - X_{\max}^{μ} will move in the same direction!
- or vice versa

Modifying CR interaction models: which way to go?

Changing the treatment of p – air interactions?

- this impacts only the initial stage of EAS development
 - further cascade development – dominated by pion-air collisions

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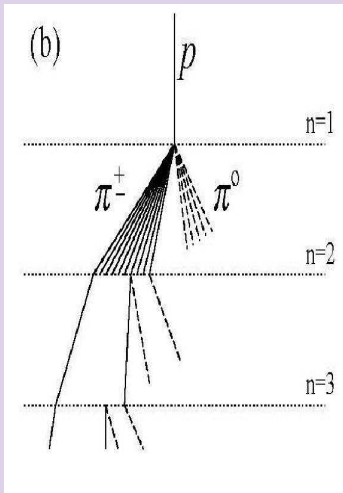
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Changing the treatment of π – air collisions ('Achilles & Tortoise')

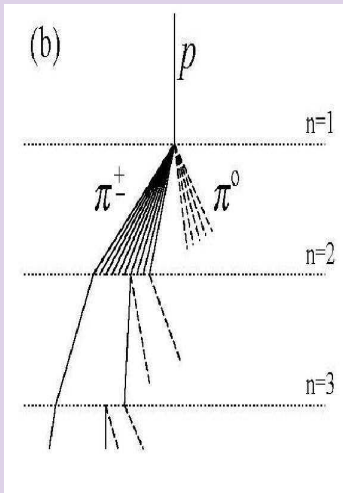
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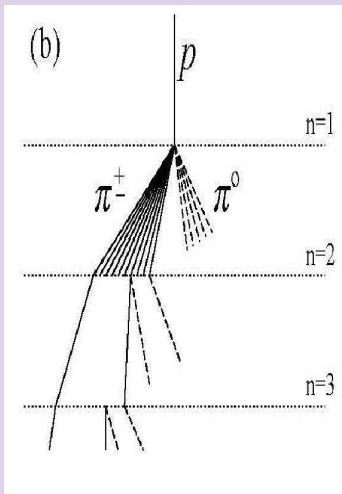
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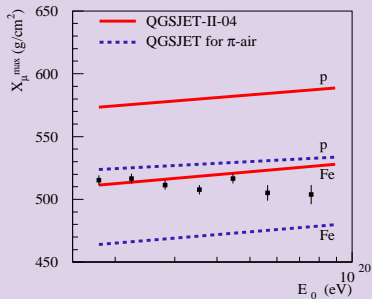
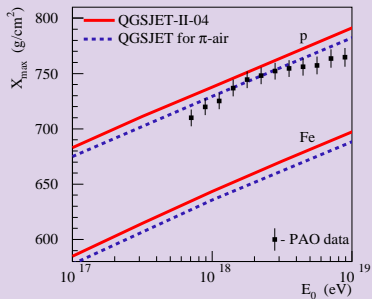
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- affects every step in the multi-step hadron cascade
 - \Rightarrow cumulative effect on X_{max}^{μ}
- but: **only the first few steps in the cascade impact X_{max}**
 - after few steps, most of energy channelled into e/m cascades
 - \Rightarrow **much weaker effect on X_{max}**



Modifying CR interaction models: which way to go?

E.g., replacing QGSJET-II by the old QGSJET, for π – air collisions

- \Rightarrow higher $\sigma_{\pi\text{-air}}^{\text{inel}}$, larger $N_{\pi\text{-air}}^{\text{ch}}$ & $K_{\pi\text{-air}}^{\text{inel}}$

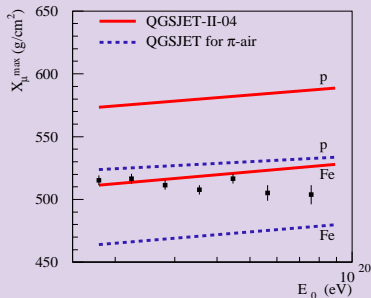
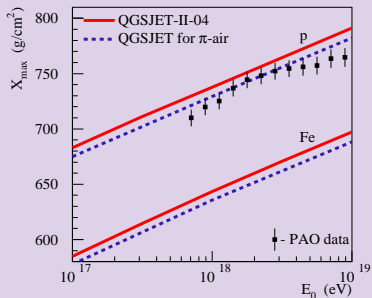


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NB: higher $\sigma_{\pi\text{-air}}^{\text{inel}}$ & $N_{\pi\text{-air}}^{\text{ch}}$ with current models – very challenging

- old QGSJET – outdated; known to overestimate particle production in π – air collisions
- needed: **drastic increase of gluon density in pions?!**